

Free Fields for Chiral 2D Dilaton Gravity. *

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Abstract

We give an explicit canonical transformation which transforms a generic chiral 2D dilaton gravity model into a free field theory.

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In the last years, a great effort has been dedicated to study two-dimensional gravity theories. They provide a simplified setting to study quantum aspects of four-dimensional gravity. The simplest one of these models was introduced by Callan, Giddings, Harvey and Strominger (CGHS) This theory can be canonically mapped into a theory of two free fields with a Minkowskian target space [2]. In the new variables it is possible to carry out a consistent functional Schrödinger quantization of the matter-coupled CGHS theory [3, 4]. This result has been recently extended to the Jackiw-Teitelboim model and the model with an exponential potential by finding explicit canonical transformations which also convert these theories into a theory of two free fields with a flat target space of Minkowskian signature [5]. In fact, it has been proven that such a transformation exists for an arbitrary model of 2D dilaton gravity [6] although the explicit form remains elusive except for the above mentioned cases. In this paper we will give an explicit expression for the canonical transformation which maps a generic 2D dilaton gravity model into a free field theory when chiral matter is coupled.

It is well-known that the action of a generic 2D dilaton gravity model can be brought to the form [7]

$$S = \int d^2x \sqrt{-g} \left[R\phi + V(\phi) + \frac{1}{2} (\nabla f)^2 \right], \quad (1)$$

by a conformal reparametrization of the fields. The equations of motion derived from (1) are

$$R + V'(\phi) = 0, \quad (2)$$

$$\nabla_\mu \nabla_\nu \phi - \frac{1}{2} g_{\mu\nu} V(\phi) = \frac{1}{2} T_{\mu\nu}^f = \frac{1}{2} \left(\nabla_\mu f \nabla_\nu f - \frac{1}{2} g_{\mu\nu} (\nabla f)^2 \right), \quad (3)$$

$$\square f = 0. \quad (4)$$

This theory is equivalent, via a canonical transformation, to a theory of three free fields. However, as we have already mentioned the explicit form of the field transformation is unknown, up to some particular cases. If we restrict the phase space of the theory by imposing the condition $T_{\mu\nu}^f T^{f\mu\nu} = 0$, which physically means to restrict the matter degrees of freedom to (let say) left-movers, the canonical transformation relating dilaton-gravity and free fields variables of the

modified (chiral) theory can be explicitly given in the general case. A convenient form for the metric, which is useful to study the chiral matter-coupled theory, is the following [8]

$$ds^2 = e^{2\rho} A dv^2 + 2e^\rho dv d\phi, \quad (5)$$

where ϕ is the dilaton field. In this way the dilaton is taken to be a coordinate and the two degrees of freedom are ρ and A . The chirality condition is automatically satisfied if we take $T_{vv}^f = T_{vv}^f(v)$ as the only non-vanishing component of the energy momentum tensor. It is easy to check that with this condition the covariant conservation of the energy momentum tensor (i.e, the equation of motion for the matter field) holds identically. The different components of the equation (3) are

$$\rho_{,\phi} = -\frac{1}{2}T_{\phi\phi}^f = 0, \quad (6)$$

$$-e^\rho A_{,v} = T_{vv}^f = \frac{1}{2}f_{,v}^2, \quad (7)$$

$$A_{,\phi} + V(\phi) = -A\rho_{,\phi}. \quad (8)$$

From equation (6) we see that ρ is a chiral field $\rho = \rho(v)$ while equation (8) can be written as

$$(A + J(\phi))_{,\phi} = 0, \quad (9)$$

where $\frac{dJ(\phi)}{d\phi} = V(\phi)$. So, the solution can be expressed as

$$A + J(\phi) = -2E, \quad (10)$$

with $E = E(v)$ an arbitrary chiral function. This function turns out to be the local energy [9]

$$E = \frac{1}{2} [(\nabla\phi)^2 - J(\phi)], \quad (11)$$

which is a conserved quantity when there is no matter at all. The constraint equation (7) becomes now

$$2e^\rho E_{,v} = T_{vv}^f. \quad (12)$$

Therefore the general solution can be written as

$$ds^2 = -e^{2\rho} \left(\int e^{-\rho} T_{vv}^f dv + J(\phi) \right) dv^2 + 2e^\rho d\phi dv, \quad (13)$$

with $\rho = \rho(v)$ an arbitrary function which is associated with reparametrizations of the coordinate v . If we choose $\rho = 0$ the above expression is a generalization of the Vaidya solution of spherically symmetric gravitational collapse for an arbitrary 2D dilaton gravity model. Our aim now is to employ this result to construct a canonical transformation which maps the theory into a free field theory. To this end we shall work out the symplectic 2-form on the phase-space constrained by the chirality condition. A useful way to achieve this is by using the covariant phase-space formalism [10] which automatically incorporates the constraint in a consistent way. The symplectic form of the theory is given by

$$\omega = \delta \int_{\Sigma} j^{\mu} d\sigma_{\mu}, \quad (14)$$

where δ stands for the exterior differential on phase space and j^{μ} is given by

$$j^{\mu} = -\phi(g^{\alpha\beta}\nabla^{\mu}\delta g_{\alpha\beta} - g^{\mu\alpha}\nabla^{\beta}\delta g_{\alpha\beta}) + \nabla^{\mu}\phi g_{\alpha\beta}\delta g^{\alpha\beta} - \nabla_{\alpha}\phi\delta g^{\alpha\mu} + \nabla^{\mu}f\delta f. \quad (15)$$

To evaluate ω on the space of chiral solutions it is convenient to choose $\phi = \text{constant}$ as the initial data surface. The result is

$$\omega = \int dv [\delta(-2E)\delta e^{\rho} + \delta f\delta f_{,v}], \quad (16)$$

implying the following Poisson bracket of dilaton-gravity variables

$$\left\{ -2E(v), e^{\rho(\tilde{v})} \right\} = \delta(v - \tilde{v}). \quad (17)$$

This is a generalization of the reasoning used in Ref[11] for the particular case of the CGHS theory.

The phase-space of the chiral theory is, therefore, made out of two arbitrary (chiral) functions $(E(v), \rho(v))$ in addition to the left-mover matter field $f(v)$.

Defining

$$X = -2E, \quad (18)$$

$$\Pi = e^{\rho}, \quad (19)$$

the constraint (12) becomes the constraint of a free field theory

$$T_{vv}^f + \Pi X_{,v} = 0. \quad (20)$$

The above formulas are the main result of this paper and provide the explicit realization of the canonical transformation proposed in Ref[6] for the chiral matter case. To finish this paper it is interesting to comment that in the absence of matter ($T_{vv}^f = 0$), E is a constant of motion and the Poisson bracket (17) converts into

$$\left\{ -2E, \int dv e^\rho \right\} = 1, \quad (21)$$

indicating that the only (global) degree of freedom is $C = -2E$ and the canonically conjugated variable is $P = \int dv e^\rho$ in agreement with [12].

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References

- [1] C. G. Callan, S. B. Giddings, J. A. Harvey and A. Strominger, *Phys. Rev. D* 45, 1005 (1992).
- [2] E. Verlinde and H. Verlinde, *Nucl. Phys. B* 406 (1993) 43.
S. Hirano, Y. Kazama and Y. Satoh, *Phys. Rev. D* 48 (1993) 1687.
A. Mikovic and M. Navarro, *Phys. Lett. B* 315 (1993) 267.
D. Cangemi, R. Jackiw and B. Zwiebach, *Ann. Phys. (N.Y.)* 245, 408 (1996).
- [3] E. Benedict, R. Jackiw and H.-J. Lee, *Phys. Rev. D* 54 (1996) 6213.
- [4] K. V. Kuchar, J. D. Romano and M. Varadarajan, *Phys. Rev. D* 55, 795 (1997).
- [5] J. Cruz, J. M. Izquierdo, D. J. Navarro and J. Navarro-Salas, "Free fields via Canonical Transformations of Matter-coupled 2D Dilaton Gravity Models" hep-th/9704168 (Phys. Rev. D, in press)
- [6] J. Cruz and J. Navarro-Salas, *Mod. Phys. Lett. A* 12, 31 (1997) 2345.

- [7] T. Banks and M. O'Loughlin, *Nucl. Phys.* B 362, 649 (1991).
- [8] J. M. Bardeen, *Phys. Rev. Lett.* 46 (1981) 382.
- [9] J. Gegenberg, G. Kunstatter and D. Louis-Martinez, *Phys. Rev.* D51 (1995) 1781.
- [10] C. Crnovic and E. Witten, in " *Three Hundred Years of Gravitation*", eds S. W. Hawking and W. Israel (Cambridge, 1987) p. 676. G. Barnich, M. Henneaux, and C. Schomblond, *Phys. Rev.* D 44, 939 (1991)
- [11] K. Schoutens, E. Verlinde and H. Verlinde, *Phys. Rev.* D 48 (1993) 2670.
- [12] D. Louis-Martinez, J. Gegenberg and G. Kunstatter, *Phys. Lett.* B 321, 193 (1994).